



# The Role of Nano-crystallites on Conduction Mechanisms of Current Through Ag Gridlines of Si Solar Cells

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## ABSTRACT

*In order to understand the impact of nano-crystallites on current transport mechanisms in screen-printed c-Si solar cells with lowly-doped emitter, Te-glass based Ag pastes with different transition temperatures ( $T_g$ ) were used. The Te-glass with lower  $T_g$  showed lower  $R_c$  than the one with higher  $T_g$  due to the formation of nano-crystallites in the glass layer. These nano-crystallites enhance the conductivity of the glass and lead to higher fill factor (FF). The nature of these nano-crystallites was first identified by the Raman spectrometry and the peaks at  $76\text{ cm}^{-1}$ ,  $119\text{ cm}^{-1}$  and  $145\text{ cm}^{-1}$  were corresponding to  $\text{Ag}_2\text{Te}$  and  $\text{PbTe}$ . The conductive-AFM further confirmed the high conductivity of these nano-crystallites without pyramidal Ag crystallites, which means the current transporting from Si emitter to Ag gridlines is mainly through the nano-crystallites in the glass.*

## INTRODUCTION

For solar electricity to reach \$0.03 kW/h in 2030 as predicted by the Department of Energy (DOE) [1], the efficiency of solar cells must approach the material limit. For instance, Si which holds the largest market share in solar industries, its theoretical efficiency can be ~30% [2]. However, since the efficiency is a product of the open circuit voltage ( $V_{oc}$ ), short circuit current density ( $J_{sc}$ ) and FF, these parameters must be increased simultaneously. In a study by ITRPV [3], the lowly-doped emitter with sheet resistance  $\sim 140\text{ }\Omega/\text{sq}$  is the target to achieve high efficiency by enhancing both  $V_{oc}$  and  $J_{sc}$ , because the lowly-doped emitter is more transparent to photons. But it is hard to contact such emitter with high  $FF > 82\%$  and it requires a new understanding of the screen-printed contacts on the lowly-doped emitter ( $\sim 140\text{ }\Omega/\text{sq}$ ).

To achieve high *FF*, the total series resistance ( $R_s$ ) encompassing emitter, gridline, busbar, bulk, back and contact (Si/Ag-gridline) should be very low. Technically, the emitter resistance is addressed by decreasing the finger spacing; gridline resistance is reduced by improving the paste rheology to fabricate continuous and thick gridlines plus increasing the number of busbars; busbar resistance is decreased by increasing the weight percent (wt.%) of Ag solids in the paste and avoiding line breakage; bulk resistance is fixed and the resistivity normally is  $\sim 2 \Omega\text{-cm}$  to reduce the light; back resistance is lowered by having uniform Al alloying in the  $\sim 10\%$  back contact area and for the aluminium back surface field (Al-BSF) solar cells, uniform BSF is important. This leaves the front contact resistance (Si/Ag-gridline), which is said to depend on the emitter peak surface concentration. The highest *FF* reported today on screen-printed solar cells, for example, Al-BSF cell is 80.92% [4] and that of the passivated emitter and rear cell (PERC) is 81.49% [5]. In this work, the role of nano-crystallites formed in the glass on current conduction between Si and gridlines is assessed. The formation of such nano-crystallites in the glass is seen to reduce the contact resistance and lead to high *FF*.

## THEORY

### Contact resistance and current transport mechanisms at the Si/Ag-gridline contacts

According to Goetzberger [6], contact resistance ( $R_c$ ) is given as

$$R_c = \frac{\rho_c}{l \times L} \quad (1)$$

where  $l$  is the length and  $L$  is the width of the gridline and  $\rho_c$  is specific contact resistance given by [7]

$$\rho_c = \frac{k}{qTA^*} \cdot \exp\left(\frac{4\pi\sqrt{\varepsilon_{Si}m^*}}{h} \cdot \frac{\Phi_{Bn}}{\sqrt{N_s}}\right) \quad (2)$$

where  $k$  is the Boltzmann constant,  $q$  is the elementary charge,  $T$  is the temperature in  $K$ ,  $A^*$  is the effective Richardson constant,  $\varepsilon_{Si}$  is the permittivity of Si,  $m^*$  is the effective mass of the charge carriers,  $h$  is the Planck constant,  $\Phi_{Bn}$  is the metal-semiconductor barrier height and  $N_s$  is the surface doping concentration.

From equation (1), the geometry ( $l$  and  $L$ ) of the gridline has an impact on  $R_c$  and could be optimized by changing the number of busbars (BB) [8] and narrowing the gridline separation [9]. Eventually,  $R_c$  is dependent on  $\rho_c$ . According to equation (2), lowly-doped emitter has higher  $\rho_c$  due to low  $N_s$  and hence higher  $R_c$ . In addition, for the screen-printed technology, the Ag gridlines are contacting not right on the top of Si surface, but several nanometers underneath it. It was found that the concentration of phosphorus decreased dramatically as the depth and it showed from the top surface to below 33 nm, the  $N_s$  decreased from  $5 \times 10^{20} \text{ cm}^{-3}$  to  $1 \times 10^{20} \text{ cm}^{-3}$  [10]. Moreover, the contact between Ag gridlines and Si emitter is not pure metal-semiconductor contact. It has a thin glass layer and metal crystallites at the interface [11]. The effect of the thin glass layer and the Ag crystallites on the current transport mechanisms is unclear. One well-accepted hypothesis is that the major current flow into the Ag gridlines is through the pyramidal Ag crystallites which directly contact with bulk Ag. This hypothesis was supported by conductive-AFM [12], microscopic I-V measurement [13] and theoretically calculation [14]. However, some researchers found that the pyramidal Ag crystallites were not necessary for a low  $R_c$  and high efficiency was achieved without pyramidal Ag crystallites [15]. In addition, the electron tunnelling assisted by nano-Ag colloids in the interface glass played a more important role in current transport mechanisms than the Ag crystallites. The more Ag colloids in the glass layer, the lower the  $R_c$  [16].

To understand the current transport mechanisms in the solar cell front contacts and reduce  $R_c$ , one main method is to modify the Ag paste constituents, which contain metal powder (~85%), glass frits (~5%), organic binder (~10%) and additives such as phosphorus dopant (PV168), Al, Cu etc. The addition of phosphorus dopant [17] did not yield any useful results because of the longer time it required for the phosphorus doping to the contact region, the higher firing temperature ( $\geq 835$  °C) and forming gas annealing [18]. In addition, the Al additives needed to be fired under optimized temperature to avoid forming Al-Si alloy and shunting the p-n junction [19]. Furthermore, the Cu additive led to the formation of  $\text{CuO}_x$  during the firing and requires forming gas annealing[20].

The alternative way to modify paste is using the glass frits with different transition temperatures ( $T_g$ ) and crystallization behaviours [11, 21]. By changing the ratio of  $\text{TeO}_2/\text{ZnO}$  [22],  $\text{TeO}_2/\text{PbO}$  [23, 24], and  $\text{TeO}_2/\text{Bi}_2\text{O}_3$  [25], the  $T_g$  of glass frits were changed. It was found that during the contact formation process, the glass frits with too low  $T_g$  started to flow earlier during the firing. There were two results: (1) the glass frits had earlier fluidization to etch the anti-reflection coating (ARC) and an earlier dissolution of Ag particles [11]. As a result, large Ag-crystallite precipitates were formed to penetrate the junction. (2) The glass frits would soften and flow more easily to form a thicker glass layer, which prevented photoelectrons from being collected. For high  $T_g$  glass frits, they needed higher sintering temperature and had insufficient time to wet the Si surface, which caused an incomplete etching of ARC and had gaps between gridlines and emitter. Thus, the *FF* was low [24]. However, the function of  $\text{TeO}_2$  in glass frits is not completely understood and the formation of  $\text{Ag}_2\text{Te}$  and  $\text{PbTe}$  after contact formation process have not been studied. In this paper, (1) Raman Spectrometer, (2) conductive-AFM, (3) SEM, and (4) EDS analyses have been used to elucidate the formation of  $\text{Ag}_2\text{Te}$  and  $\text{PbTe}$ , which are believed to decrease the  $R_c$  for lowly-doped emitter.

## EXPERIMENT

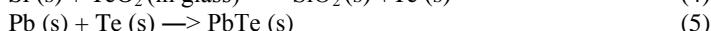
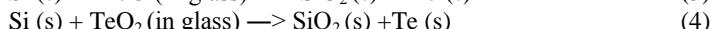
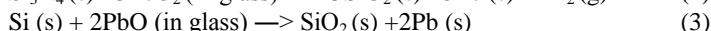
The p-type Czocralski wafers with bulk resistivity of ~2.5 Ω-cm were textured and followed by phosphorus diffusion at 890°C to form 95 Ω/sq emitter. After that, the wafer edges were isolated and the phosphorus glass was removed followed by PECVD  $\text{SiN}_x$  (73 nm) deposition. The wafers were divided into four groups for four Ag pastes (A, B, C, D) based on  $\text{TeO}_2$  glass with different  $T_g$ . Three-busbar structure for full Al-BSF was used in this experiment. The four front Ag pastes were screen-printed in turn onto the wafers with back Al contact already printed and dried. After drying the front Ag paste, the cells were co-fired in the rapid thermal processing (RTP) infrared belt furnace at 230 inch per minute (ipm) at 815°C peak temperature. The fabricated cells were first characterized by light I-V measurements, then one cell from each group was cut into 2 mm strip and the contact resistance was measured.

To investigate the contact interface between the underlying Si and gridline, the cut samples were generated into three sets, (a) as fired (AF); (b) Ag metal removed by  $\text{HNO}_3$  but the glass remained ( $\text{HNO}_3$ ); and (c) the remained glass removed by  $\text{HF}$  (HF). However, from the contact resistance measurements, the contact resistance for pastes B to D were similar, therefore, other analyses were carried out on samples only with pastes A and D, with highest and lowest  $R_c$  and  $R_s$ . The microstructure, elemental composition and conductive properties of these samples were evaluated with field-emission scanning electron microscope (FESEM) with EDS (FEI Verios 460L), Raman spectrometer (HORIBA, Xplora™PLUS) excited at 532 nm and conductive-AFM (Asylum MFP-3D).

## RESULTS AND DISCUSSION

### Chemical reactions during contact co-firing step

The front Ag contacts are formed by sintering Ag paste under IR belt with temperatures ranging from 400-815°C for a short time. As the temperature increases from low to high, (a) the organic binder burns out; (b) glass frits start to melt and diffuse towards the wafer surface; (c) the melted glass etches away the ARC layer on the wafer surface through the reduction reactions (1) and (2); after ARC removal, the glass starts to react with the Si substrate through reactions (3) [26] and (4) [27]; meanwhile, the formed Pb and Te could further react with Ag and form PbTe and Ag<sub>2</sub>Te from reactions (5) and (6).



Based on Gibbs free energy, the thermodynamic potential for each possible reaction is shown in Figure 1(a) where the potential of reactions (5) and (6) at 850°C is respectively -69.54 kJ/mol and -59.36 kJ/mol. It is obvious that for the reactions (1) ~ (6), the Gibbs free energy is always negative from 300°C to 850°C, which means each reaction can happen spontaneously.

### $R_c$ , $R_s$ and FF of samples with four pastes

Figure 1b and 1c depict  $R_c$ ,  $R_s$  and the corresponding FF for the four pastes. Paste A shows the highest  $R_c$  and  $R_s$  with lowest FF. The high  $R_c$  for contacts with paste A can be attributed to the thicker glass layer (~3 μm) in Figure 2a at the interface of Si/Ag-gridline, while the lower  $R_c$  for paste D contacts is evident in the thinner glass layer of ~0.7 μm in Figure 2b.

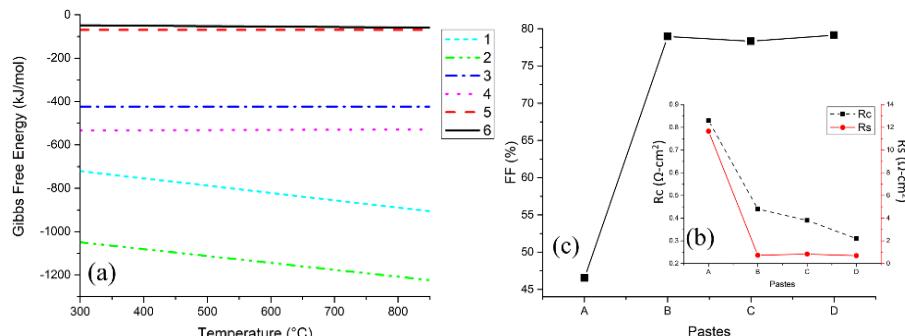


Figure 1, (a) Gibbs free energy for possible reactions (1) ~ (6) under temperature from 300°C to 850°C. (b)  $R_c$  and  $R_s$  of pastes A-D with (c) FF

## SEM and EDS analyses

To understand the metal-semiconductor contact interface for the samples with pastes A and D, the microstructure analyse was carried out as shown in Figures 2a and 2b, respectively. The contact with paste A exhibits a thick non-uniform interface glass layer of  $\sim 3\mu\text{m}$  thickness as opposed to those of D with a uniform glass layer of  $\sim 0.7\mu\text{m}$  and some Ag nano-particles ( $\sim 50\text{ nm}$ ) inside. Figures 2c and 2d, respectively, show the corresponding EDS for the glass frits in pastes A and D along with the Pb:Te ratio of 3.1 and 1.6 respectively. The Pb:Te ratio is known to impact the  $T_g$  of the glass [11, 21]. Since the glass frit in paste A had higher  $T_g$  than that of D, it is possible that higher firing temperature is needed for the glass frits to melt, wet the Si surface, uniformly etch the  $\text{SiN}_x$ , and then form a thin glass layer with nano-particles inside as in D.

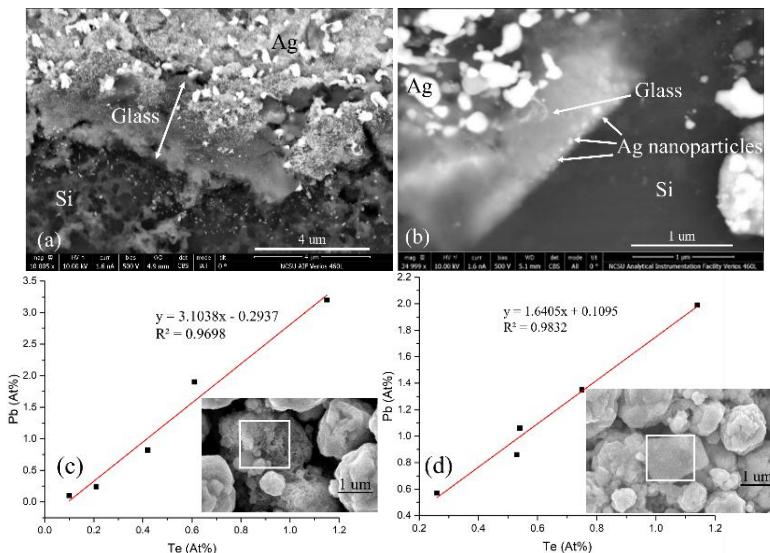


Figure 2, SEM images of samples with pastes (a)A and (b)D with Pb:Te ratio of the glass frits in pastes (c) A and (d) D.

Figure 3a and 3b show the EDS and SEM of the sample with paste A after the Ag gridline was removed with  $\text{HNO}_3$ . This is an attempt to investigate the location of the nano-particles in the glass, as  $\text{HNO}_3$  only removes the Ag gridline without attacking the glass. As seen in Figure 3b, there are only few nano-crystallites and the elemental analysis shows no Te in the analysed region. On the contrary, Figure 3d for paste D, shows micro-sized bright areas which are large nano-crystallites and such nano-crystallites contain Pb, Ag and Te.

In order to further ascertain the location of these large nano-crystallites, the glass layer was etched off with  $\text{HF}$ . Because the sample with paste A did not show any crystallites, EDS was carried out only on paste D after  $\text{HF}$  treatment. As shown in Figure (f), there was few Ag crystallites in the Si for D which indicates the large nano-crystallites were formed right in the glass layer.

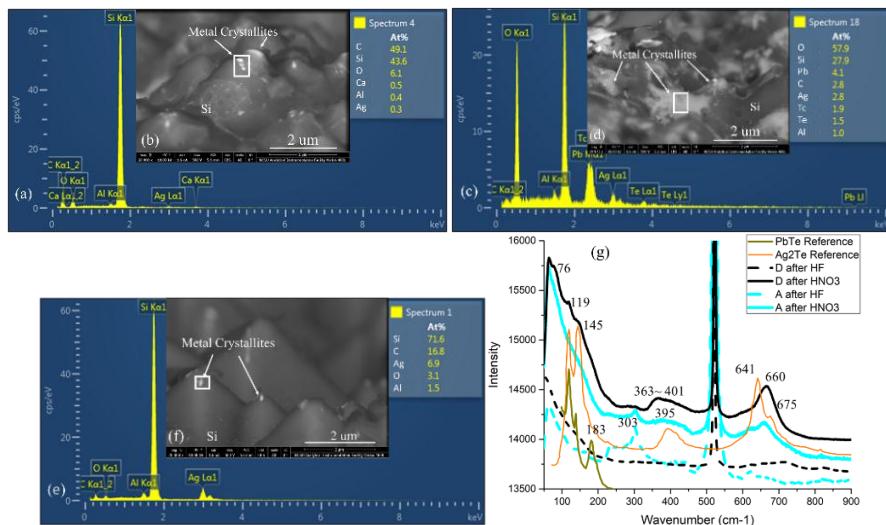


Figure 3. EDS elemental analyses of pastes (a) A, (c)D after  $HNO_3$  treatment and (e) D after HF, where the analysed areas are marked by boxes as shown in SEM images: (b) A, (d) D after  $HNO_3$  and (f) D after HF. And Raman spectroscopy (g) of pastes A and D after  $HNO_3$  and HF treatment.

### Raman spectrometer

In order to fully understand the nature of the nano-crystallites formed in the glass layer, the Raman spectroscopy was used. As shown in Figure 3(g), after the  $HNO_3$ , D exhibits peaks at  $76\text{ cm}^{-1}$ ,  $119\text{ cm}^{-1}$  and  $145\text{ cm}^{-1}$  which match the  $Ag_2Te$  and  $PbTe$  [28-30]. The peaks around  $360\text{--}400\text{ cm}^{-1}$  and  $660\text{ cm}^{-1}$  are binding vibrations of  $TeO_2$  [31]. For A, it only has  $TeO_2$  peaks and no  $Ag_2Te$  and  $PbTe$ . After the HF treatment, both pastes A and D only show Si peaks at  $303\text{ cm}^{-1}$  and  $520\text{ cm}^{-1}$ . The Raman spectroscopy confirms that the nano-crystallites formed in the glass for D contains  $Ag_2Te$  and  $PbTe$ .

### Conductive-AFM

In order to understand the distribution of nano-crystallites in the current transport, the glass conductivity was studied by conductive-AFM, where the samples had Ag gridlines removed by  $HNO_3$ . The current was measured on the front side after applying a 10V bias on the backside of the sample. In Figure 4d, the sample with paste D has micro-sized areas which is very conductive with current over 10 nA, while most area of sample A is non-conducting. The conductive area in D matches the size and shape of the large nano-crystallites in Figure 3(d). This suggests that the nano-crystallites in the glass layer enhance the glass conductivity and without these nano-crystallites, the glass is insulator. Thus, the current transports from Si emitter to Ag gridlines mainly through these nano-crystallites of  $Ag_2Te$  and  $PbTe$ .

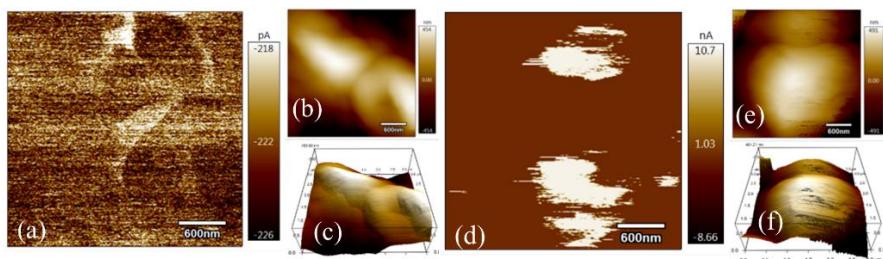


Figure 4, Conductive-AFM measurements on samples with pastes A (a~c) and D (d~f) according to current (a&d), height (b&e) and 3D geometry (c&f) measurement.

## CONCLUSIONS

In this study, it is found that the Pb to Te ratio as well as the  $T_g$  of the glass frits can impact the reformed glass at interface of Si/Ag-gridline after contact sintering step. The glass with lower  $T_g$  enhances the uniform wetting of the Si surface and  $\text{SiN}_x$  etching as well as the formation of thinner glass layer with large nano-crystallites. These large nano-crystallites embedded in the glass contain  $\text{Ag}_2\text{Te}$  and  $\text{PbTe}$  and increase the glass conductivity. Thus, the current transport from the Si material to Ag gridline is mainly through these nano-crystallites in the glass. In addition, since the pyramidal Ag crystallites were not found in the emitter after the glass layers removal, it suggests that the interface glass plays a more important role in current transport mechanisms than the Ag crystallites.

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