

①

IN-class  
NOTES  
3/29/23

EXAMPLE  $N$ -monatomic molecules in Heated Box  
Ref: 4.4 Hill's

1)  $R_{\text{IND}} \text{ MODE} \Rightarrow$  NO INTERNAL D.O.F.

2)  $R_{\text{IND}} \text{ MODE} \Rightarrow$  Atoms; taking into account nuclear

Q. Dynamics  $\oplus$   
e. electron dynamics  
~~atomic~~ about nuclei

$\oplus$   
translational Quantum energy of Atom

C.O.M. we've done this

B) 1)  $H_{\text{ATOM}} = H_{\text{nucleus}} + H_{\text{el.}} + H_{\text{translation C.O.M.}}$

a) RANDOM dynamics of 3 modes statistical ~~indep.~~ indep.

C) Temp of Box/gas (2)  
 is high enough  
 and/or  $N/V$  small  
 enough that  
Quantum stat mechs  
~~can~~ is app'd by  
classical limit

D) (2,7)

For  $cc$ :  $Q_{\text{QUANTUM}} = \frac{Q_{\text{classical}}}{N!}$  high T / low den  
 $Q_{\text{classical}} = (Q_{\text{atom}})^N$   
 $Q_{\text{classical}} = \frac{(Q_{\text{atom}})^N}{N!}$

⊕  $N$  exponent  $\rightarrow$  Assump 2?  $\rightarrow$  indep dynamics  
 of each atom  
 is stat. indep

⊕  $N!$   $\Rightarrow$  proven in upcoming lecture

- ⊕ Mill
- ⊕ GIBB'S correction

B)  $\Psi_{\text{atom}}$

a) since (3.1)

$$H_{\text{atom}} = H_{\text{nucleus}} + H_{\text{elec}} + H_{\text{trans}}$$

b) + then (3.2)

$$\hat{H}_{\text{atom}} | \Psi_{\text{atom}} \rangle = E_{\text{atom}} | \Psi_{\text{atom}} \rangle$$

c) ~~but~~ (3.1) allows factoring (3.2) =

S.S. nucleus

(c1)

$$\hat{H}_{\text{nuc}} | \Psi_{\text{nuc}} \rangle = E_{\text{nuc}} | \Psi_{\text{nuc}} \rangle$$

??  
nucleus

S.S. electron

(c2)

$$\hat{H}_{\text{elec}} | \Psi_{\text{elec}} \rangle = E_{\text{elec}} | \Psi_{\text{elec}} \rangle$$

electron

C.C.M.

(c3)

$$\hat{H}_{\text{trans}} | \Psi_{\text{trans}} \rangle = E_{\text{trans}} | \Psi_{\text{trans}} \rangle$$

d)

(c1) $\Rightarrow$	$E_{n_1}, E_{n_2}, \dots$
(c2) $\Rightarrow$	$E_{e_1}, E_{e_2}, \dots$
(c3) $\Rightarrow$	$E_{t_1}, E_{t_2}, \dots$

④

F) 
$$\left[ \epsilon_{\text{Atom}} = \epsilon_{n_i} + \epsilon_{e_j} + \epsilon_{+k} \right] \quad (4.1)$$

M) ALSO, BECAUSE WITH ASSING  
STAT. INDEP DYNAMICS

$$q_{\text{Atom}} = q_{\text{Ave}} q_{\text{elec}} q_{\text{tran}}$$

$$\left[ q_n = q_n q_e q_t \right] \quad (4.2)$$

III  $[q_n]$  = partition for nuclear  
dynamics?

A) 
$$\left[ q_n = \sum_{s=1}^{\infty} e^{-\beta \epsilon_{n_s}} \right] \quad (4.3)$$

OR WITH ④  $W_{n_s} \equiv$  no. of distinct  
nuclear quantum  
states having  
Q. energy  $[E_{n_s}]$

A) Min:

~~NSD~~  $\left\{ \begin{array}{l} \text{preliminary guess:} \\ \text{ground state energy} \\ \text{for a ptc} = \\ \text{ptcl en. @ } T = 0K \end{array} \right.$

B) Hill:

ii)  $\boxed{\epsilon_{n_i} \approx 0}$

$\oplus$  now's mod/

iii)  $S(0K) =$

$k_B \ln \Omega$

$\downarrow$   
 $\boxed{\omega_{n_i}}$

ii)  $\boxed{\omega_{n_i} = \text{not } 0}$

$\Downarrow$   
 $\boxed{S_{\text{vib}}(0K) \neq 0}$

As, vs perfect crystals

have 0 degeneracy

@  $0K \Leftrightarrow$  there's only

$\downarrow$   $\boxed{\text{unique \& state}}$

@  $0K \left\{ \begin{array}{l} (T=0K) \end{array} \right.$

$\oplus S = k_B \ln \Omega_{\text{as}} = \underline{0}$   
crystal <sub>3</sub>

Nuclei cont'd (M.11)

⑥

⊕  $E_{n_2} = O(10^6 \text{ eV})$

⊕  $kT \approx 10^4 \text{ K}$

⊕ To get nuclei excited from  $E_{n_1} (\approx 0)$  to  $E_{n_2}$  thermally / collisionally  
 reqs  $T \sim O(10^{11} \text{ K})!!$

∴ a) 
$$g_n = \sum_{s=1} W_{n_s} e^{-\beta E_{n_s}}$$

b) 
$$g_{\text{nuclear}} = W_{n_1} e^{-\beta E_{n_1}} + W_{n_2} e^{-\beta E_{n_2}} + \dots$$

$\beta E_{n_2} \sim \frac{k_B(10^{11} \text{ K})}{k_B(10^4 \text{ K})} e^{-10^7}$

c) For temp  $\approx 10^4 \text{ K}$   
 $P_n(E_{n_2}) \Rightarrow P_{\text{Boltzmann}}(E_{n_2})$

where  $P_{\text{bottom}}(\epsilon_{n_2}) = \frac{e^{-\beta \epsilon_{n_2}}}{\mathcal{E}_{\text{fuel}}}$  ⑦

④  $= \frac{e^{-10^{-7}}}{w_{n_1}}$

$P(\epsilon_{n_2}) \rightarrow 0$

⑤  $P_{\text{Boltz}}(\epsilon_{n_1}) = \frac{e^{-\beta \epsilon_{n_1}}}{\sum e^{-\beta \epsilon_{n_i}}} = 1$

⑥ Temps  $\Rightarrow \ll 10^{10}$  K  
all atomic nuclei  
are in their  
ground state

# IV) [Ge]

Alkali  $\oplus$  physical factors  $\oplus$

Alkali Atoms K, Li, etc

$$\Delta E_e = E_{e2} - E_{e1}$$

$\uparrow$  first exci       $\uparrow$  ground state elec. cloud ch.

$$\Delta E_e \sim 1.5 \text{ eV} \approx 15000 \text{ K}$$

Halogen Ar, Kr, Xe, Ne

$$\Delta E = 10 - 20 \text{ eV} \leftarrow 10^5 - 2(10^5) \text{ K}$$

Halogens F, Cl, etc

$$\Delta E = 0.05 \text{ eV} - 1 \text{ eV}$$

$\updownarrow$  10<sup>5</sup>00K       $\updownarrow$  10<sup>4</sup>K



4)  $g_n = w_n$

ASIDE on degeneracy: when a multiple ~~of~~ distinct states have same energy  
microstate system = prob can

of example: see

translational



$$g = e^{-\beta \epsilon_0} + e^{-\beta \epsilon_1} + \dots$$
$$= [e^{-\beta \epsilon_0} + e^{-\beta \epsilon_1} + \dots] + [e^{-\beta \epsilon_2} + \dots]$$

1)  $g_{states} = \sum_{s=0} e^{-\beta \epsilon_s}$

2)  $g_{levels} = \sum w(\epsilon_s) e^{-\beta \epsilon_s}$

3)  $w(\epsilon_s) = \text{deg of } \epsilon_s$

4)  $p(\epsilon_s)_{states} = \frac{e^{-\beta \epsilon_s}}{g}$

5)  $p_{levels}(\epsilon_s) = \frac{w(\epsilon_s) e^{-\beta \epsilon_s}}{\sum_{s'} w(\epsilon_{s'}) e^{-\beta \epsilon_{s'}}$

for small  $\Delta \epsilon$ 's

6)  $\sum_{s'} w(\epsilon_{s'}) e^{-\beta \epsilon_{s'}} = w(\epsilon_1) e^{-\beta \epsilon_1} \Delta \epsilon_M$   
 $+ w(\epsilon_2) e^{-\beta \epsilon_2} \Delta \epsilon_M$

from calc  $\sum_{s'} w(\epsilon_{s'}) e^{-\beta \epsilon_{s'}} \Delta \epsilon \Rightarrow \sum = \int w(\epsilon) e^{-\beta \epsilon} d\epsilon$

$$Q_{rot} \frac{\Delta E_{rot}}{M} = [w(\epsilon_1) e^{-\beta \epsilon_1} + w(\epsilon_2) e^{-\beta \epsilon_2} + \dots] \left( \frac{\epsilon_{max} - \epsilon_{min}}{M} \right)$$

$$Q_{rot} = \left[ \frac{w(\epsilon_1) e^{-\beta \epsilon_1}}{\Delta E_{rot}/M} + \frac{w(\epsilon_2) e^{-\beta \epsilon_2}}{\Delta E_{rot}/M} + \dots \right] \frac{\Delta E_{rot}}{M} \Delta \epsilon$$

$$= \sum_i \left[ \frac{w(\epsilon_i) e^{-\beta \epsilon_i}}{\Delta E} \right] \Delta \epsilon$$

$$= \sum_i \tilde{w}(\epsilon_i) e^{-\beta \epsilon_i} \Delta \epsilon = \int_0^{\infty} \tilde{w}(\epsilon) e^{-\beta \epsilon} d\epsilon$$

$\tilde{w}(\epsilon_i) = \frac{w(\epsilon_i)}{\Delta E} = \frac{\text{no. of molecular states w/ energy } \epsilon_i}{(\text{unit energy interval})}$

$$= \int \tilde{w}(\epsilon) e^{-\beta \epsilon} d\epsilon$$

$$= \Phi =$$

$$\Phi(\epsilon) =$$

diatomic molecule

des by  $^{2S+1}L_J$

$$2P_{3/2}$$

$\Rightarrow$

$$\tilde{S} = 1/2$$

$$L = 1$$

$$J = \tilde{S} + L = 3/2$$

$$\rightarrow g = 4$$

form

term

symbol

$$2P_{1/2}$$

$\Rightarrow$

$$\tilde{S} = 1/2$$

$$L = 1$$

$$J = \tilde{S} + L =$$

$g \in \mathbb{N}$  integer between  
 $+J$  and  $-J$

⊕

$$2P_{3/2} \Rightarrow g = 4$$

$$\text{TOTAL } g = 2J + 1$$

$$\therefore \begin{array}{l} 2P_{3/2} \rightarrow g = 2\left(\frac{3}{2}\right) + 1 \\ \quad \quad \quad = 4 \\ 2P_{1/2} \rightarrow g = 2\left(\frac{1}{2}\right) + 1 \end{array}$$

$$Z_{\text{norm}} = Z_e Z_c Z_n$$

prop  
N<sub>3</sub> 10<sup>23</sup>  
v. 2  
10<sup>6</sup> eV

⊙

$Z_n$ : (a)  $E_{\text{ion}} - E_n \sim 10^{10} \text{ K}$

⇓  
⊙ atomic nuclei in many ground state

⊙  $Z_n = \sum_{s, n} w_{s, n} e^{-\beta E_{s, n}}$   
 $= \Omega_{n, 1}$

choose  
~~define~~  
 $E_{n, 1} = 0$

$Z_e$

$\epsilon_2 \approx \epsilon_1 \approx k_B T$ ; ;  $\epsilon_3 \gg k_B T$

$$P(\epsilon_2) \approx \frac{w_2 e^{-\beta \epsilon_2}}{w_1 e^{-\beta \epsilon_1} + w_2 e^{-\beta \epsilon_2}}$$

(1)

= Halogens

F, Cl, Br, I:

$$\Delta \epsilon_i = \epsilon_2 - \epsilon_i \quad \downarrow \quad \downarrow \quad \downarrow$$

0.05    0.11    0.46    0.94 eV  
 $\sim 500K \sim 100K \sim 5000K \sim 10^4K$

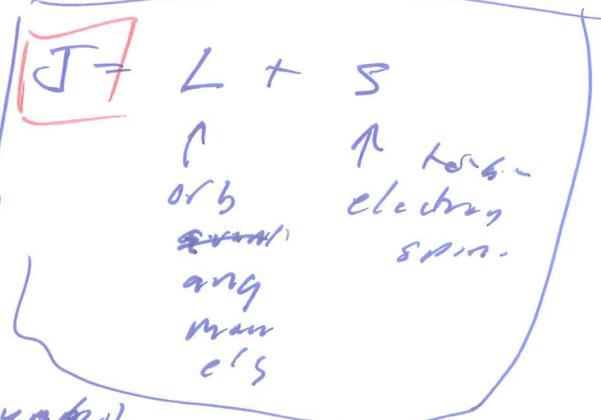
$$w_{\epsilon_1} = 4, \quad w_{\epsilon_2} = 2$$

Set  $(2S+1 L_J)$  term symbols

~~g~~

$$2S+1 = 2 \Rightarrow S = \frac{1}{2}$$

$$L = 0 \Rightarrow J = \frac{1}{2}$$



1 +

L	Symbol
0	S
1	P
2	D
3	

$$g = 2J + 1$$

P(2)

II) a)  $Q$  or  $Q$  in terms of discrete energy levels  $\epsilon_n$  (2)

$$Q = \sum_s e^{-\beta \epsilon_s}$$

b)  $Q$  in terms of molecular energy levels and degeneracy  $w$

$$Q = \sum_s w_s e^{-\beta \epsilon_s}$$

nuclei  
or  
elect  
or  
Trans

$w_s$

$w_s =$  no of q. states having energy  $= \epsilon_s$

or ~~no~~  $w_s \epsilon_s$

when  $\epsilon_s$ 's are close together  $Q$  or  $Q$  can be expressed as integral.

c)

For  $k_B T \ll \epsilon_2 - \epsilon_1$

(3)

$$Z_{cl} = \sum_{s=1}^2 \int d^3x e^{-\beta \epsilon_s} = e^{-\beta \epsilon_1} + e^{-\beta \epsilon_2}$$

class

$$\epsilon = \frac{2\pi^2 \hbar^2}{8mL^2} (n_x^2 + n_y^2 + n_z^2)$$

$$\epsilon =$$

$$a) \sqrt{n_x^2 + n_y^2 + n_z^2} = R$$

$$R = \sqrt{\epsilon_{xyz} \left( \frac{8mL^2}{\hbar^2} \right)^{1/2}}$$

$\frac{\hbar^2 2\pi^2}{mL^2}$   
 $= \hbar^2 \frac{2\pi^2}{8mL^2}$

b)  $\Phi(\epsilon_{xyz} \in k_B T) = \text{vol.}$

$$= \frac{4}{3} \pi \frac{R^3}{8}$$

$$= \left( \frac{\pi}{6} \right)^{3/2} \left( \frac{8mL^2}{\hbar^2} \right)^{3/2} (k_B T)^{3/2}$$

$\Downarrow$   
 $\epsilon$

c)

$$d) i) Q_T = \sum_{i=1} w(\epsilon_{Ti}) e^{-\beta \epsilon_{Ti}}$$

$$\oplus \epsilon_{Ti} \in \mathbb{R}_{xyz} = x, y, z$$

$$ii) Q_T = \sum_i \frac{w(\epsilon_{Ti})}{\Delta \epsilon} e^{-\beta \epsilon_{Ti}} \Delta \epsilon$$

$$\Delta \epsilon = \frac{\epsilon_{Tmax} - \epsilon_{Tmin}}{N}$$

$$= \sum \tilde{w}(\epsilon_T) e^{-\beta \epsilon_T} \Delta \epsilon$$

$$\oplus \tilde{w}(\epsilon_T) \stackrel{\text{no. of distinct states having } \epsilon = \epsilon_T}{\text{(per unit } \Delta \text{ in interval } \epsilon)}$$

$$\therefore Q_T = \int_0^{\infty} \tilde{w}(\epsilon_T) e^{-\beta \epsilon_T} d\epsilon_T$$

$$e) \tilde{w}(\epsilon_T) : a) \Phi \approx \frac{1}{6} \left( \frac{2\pi m \epsilon_T}{h^2} \right)^{3/2} \left( \frac{2\pi m \epsilon_T}{h^2} \right)^{3/2} \left( \frac{2\pi m \epsilon_T}{h^2} \right)^{3/2}$$

= no of trans & states less than  $\epsilon \leq \epsilon_T$

$$b) \tilde{w} = \frac{d\phi(\epsilon)}{d\epsilon} = \left( \frac{\text{no. of states having } \epsilon = 0}{d\epsilon} \right) \quad (5)$$

$$\text{or } \tilde{w}(\epsilon) = \frac{d}{d\epsilon} \left[ \left( \frac{8m\sqrt{2/3}}{h^2} \right)^{3/2} \epsilon^{3/2} \right]$$

$$= \frac{1}{2} \left( \dots \right)^{3/2} \frac{1}{\epsilon^{1/2}}$$

$$c) \therefore Z_+ (V, T) = \frac{1}{2} \left( \dots \right)^{3/2} \int_0^{\infty} \epsilon^{-1/2} e^{-\beta\epsilon} d\epsilon$$

$$b) \tilde{w}(\epsilon) = \frac{d\phi(\epsilon)}{d\epsilon} = \frac{d}{d\epsilon} \left[ \left( \frac{8m}{h^2} \right)^{3/2} (V) \left( \frac{\pi}{6} \right) \epsilon^{3/2} \right]$$

$$= \left[ \left( \frac{3}{2} \right) \left( \frac{8m}{h^2} \right)^{3/2} \left( \frac{\pi}{6} \right) \right] \epsilon^{1/2}$$

$$c) \therefore Z_+ (V, T) = \left[ \dots \right] \int_0^{\infty} \epsilon^{1/2} e^{-\beta\epsilon} d\epsilon$$